

Research Paper

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A germanium telluride switch optically controlled with infrared showing over 30k cycles of operation and 32-fs figure-of-merit

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Abstract

This article presents germanium telluride (GeTe)-based switches for radiofrequency (RF) applications, capable of reversible switching between their ON and OFF states through optical activation by irradiation. Unlike previous studies, the transition is induced by infrared laser pulses at a wavelength of $\lambda = 915$ nm, which is highly promising for future integration of laser sources and the proposal of fully integrated optical activation of phase change material (PCM) switches. This represents a novel approach compared to the existing literature, which primarily focuses on the ultra-violet spectrum, less suitable for on-chip optical integration. Our work also provides combined optical and thermal simulations to elucidate the challenges associated with actuating small PCM switches and demonstrates the effectiveness of PCMs at this wavelength. The study achieves bistable switching at high frequencies up to 40 GHz, with a figure of merit of 31.5 fs, despite the low GeTe conductivity of only $1.85 \cdot 10^5$ S/m. Additionally, significant advancements over the literature have been made by surpassing 30,000 cycles with optical actuation.

Introduction

Phase change materials (PCM) have a long history of application, mainly in the field of non-volatile memories [1] and the optical storage industry [2]. Beyond these traditional applications, PCMs have also found a variety of uses in integrated photonics [3]. Recently, chalcogenide PCMs, particularly germanium telluride (GeTe) and germanium-antimony-telluride, have emerged as promising materials for the design of next-generation radiofrequency (RF) switches. These switches exploit the unique properties of PCMs, such as their ability to phase shift between amorphous and crystalline states, enabling rapid and reversible transitions with significant effects on the electrical characteristics of the device. RF switches based on PCMs offer several notable advantages. First, they can handle high power levels while maintaining a low resistance-capacitance product ($R_{ON}C_{OFF}$), which is crucial for the performance of RF circuits. Second, these switches are non-volatile, meaning they retain their switching state even in the absence of power supply, thereby reducing power consumption and increasing system reliability [4]. Third, their low power consumption makes them particularly suitable for applications where energy efficiency is essential. These characteristics make PCM-based RF switches indispensable for modern RF technologies, which must meet very interesting for stringent performance and efficiency requirements. For instance, in millimeter-wave radar systems, the ability of switches to handle high frequencies with minimal losses is crucial for improving radar resolution and range. Similarly, 6G communication networks, aiming to deliver ultra-fast data transmission speeds and extremely low latency, can benefit from the integration of PCM switches to optimize RF circuit performance [5–7].

However, the activation of these PCM-based RF switches presents significant challenges. Conventional methods often rely on indirect heating using micro-heaters, which generate heat through the Joule effect [8]. This can lead to considerable parasitic capacitance, thereby limiting RF performance. Additionally, the non-uniform temperature distribution between the micro-heater and the PCM element can adversely affect the RF performance [9]. Incomplete crystallization or amorphization of the PCM can result in increased R_{ON} or C_{OFF} , thus degrading the switch cutoff frequency. Alternative solutions, such as optical actuation, have been explored [10–12]. These methods leverage the strong optical absorption of GeTe at specific wavelengths, allowing the conversion of optical energy into heat to switch the PCM for various applications. For instance, optically controlled GeTe has been used to reconfigure band-stop filters and polarizers in the THz frequency range [12]. A reconfigurable antenna utilizing GeTe deposits to

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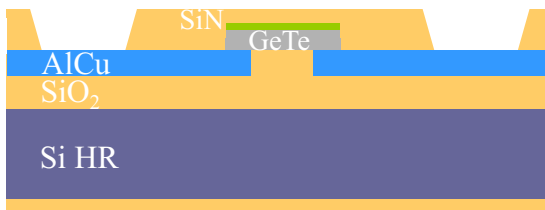


Figure 1. Switch stacking structure.

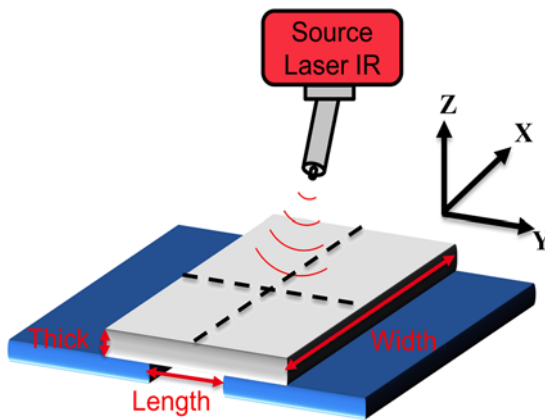


Figure 2. 3D simulation setup.

connect or disconnect radiating elements is presented in [10]. Furthermore, [11] demonstrates a large-scale RF-PCM switch, with performance characterized by frequency. Nevertheless, these approaches have been hampered by the use of expensive UV lasers that are not compatible with current CMOS integrated photonics technology. In contrast, a wavelength of 915 nm, which is the standard for integrated photonic circuits, offers a promising alternative [13, 14]. This wavelength allows for the use of low-cost, chip-scale pulsed lasers that can be easily integrated into a photonic chip.

Therefore, studying the actuation of RF PCM switches at 915 nm offers a compelling solution to the aforementioned challenges and paves the way for PCM/photonic co-integration. An earlier version of this paper was presented at EuMW 2024 and was published in its Proceedings [24]. In this context, this work presents, to the best of our knowledge, the first combined thermo-optical simulations of GeTe RF switches to highlight the issues in actuating small switches, as well as to study their optical actuation performance at an IR wavelength of 915 nm and the confirmation of those results by measurements. It should be noted that power handling measurements have already been performed on similar switches in previous work reported in [23]. The first part of this work describes the technological stack of our GeTe RF switches and the results of combined optical and thermal simulations. The second part details the conditions for optical actuation at 915 nm and the measured results that include DC and RF performances and endurance testing. The final part provides a comparison with the state-of-the-art.

Optical actuation of GeTe material

Technology and test bench description

In this study, RF switches in a Single Pole Single Throw (SPST) configuration are integrated on a coplanar waveguide utilizing GeTe, with a Ge[50%]:Te[50%] target. The GeTe is deposited at a low temperature and subsequently annealed to achieve an initial crystalline state. The switch stacking structure is shown in Figure 1. The developed process flow ensures compatibility with standard CMOS back-end-of-line processes and meets the performance requirements of PCM switches. Using this technological framework, we have developed a 3D simulation setup for the optical actuation of our switches, as illustrated in Figure 2. In this setup, two RF lines are connected by a GeTe segment, parameterized by its width (X), length (RF gap) (Y), and thickness (Z). An optical wave source with a Gaussian profile is positioned 10 μm above the GeTe to induce phase transitions between the crystalline (ON) and amorphous (OFF) states. This transition is accomplished by heating the GeTe

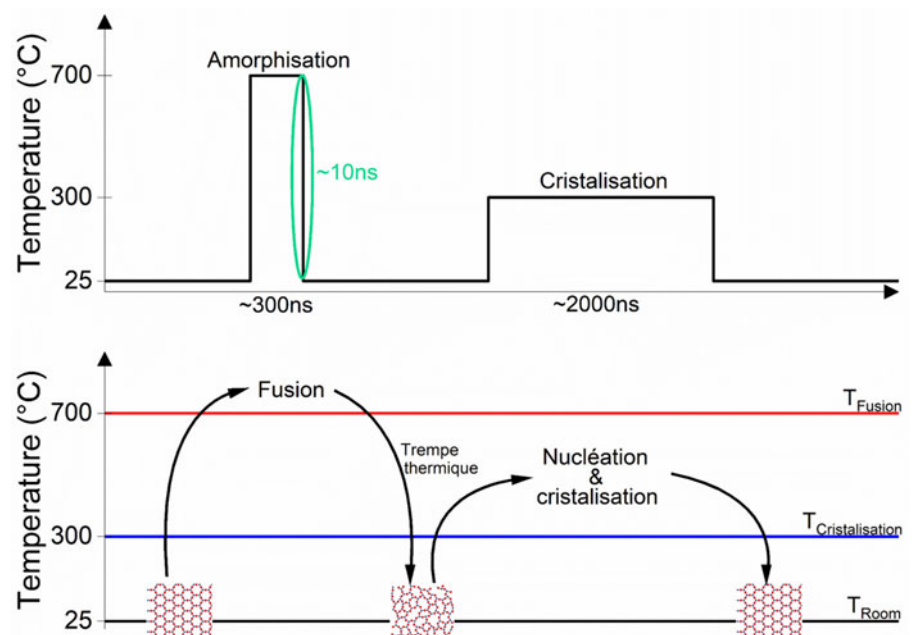


Figure 3. Reversible switching of a phase-change material using an optical pulse.

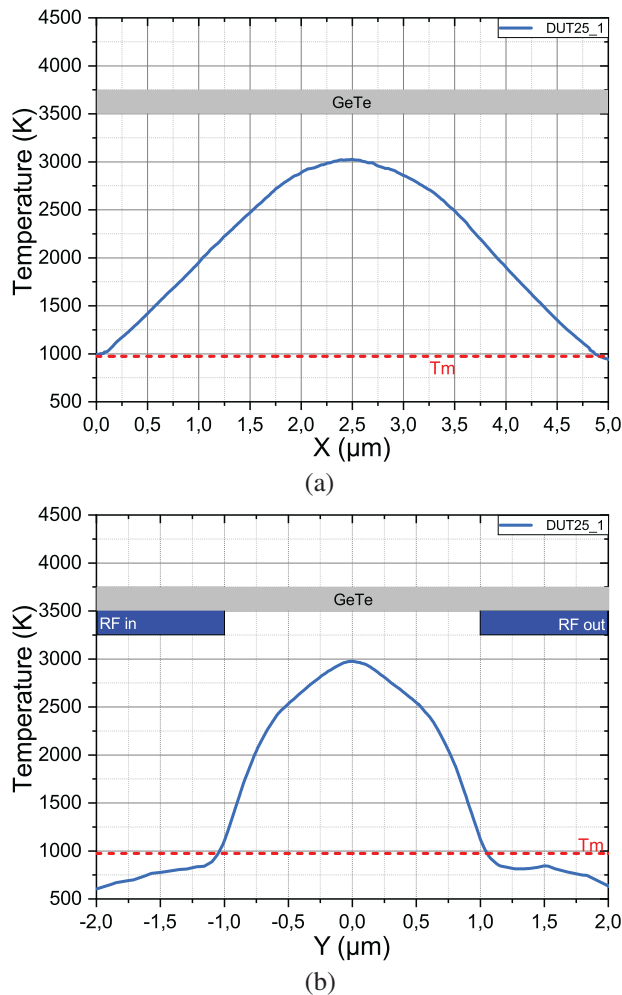


Figure 4. Temperature variation according to: (a) PCM length and (b) PCM width.

above its melting temperature ($T_m \sim 973$ K) for a brief period, followed by rapid quenching to solidify the atoms in the amorphous state, as shown in Figure 3. Crystallization is achieved by applying a pulse of longer duration (2 μ s), which heats the PCM above its recrystallization temperature ($T_c \sim 573$ K), thereby activating crystal grain growth and nucleation. This process changes the resistivity of the PCM from the amorphous (high resistivity) to the crystalline (low resistivity) state.

Optical and thermal coupled simulation

Optical and thermal simulations were conducted using ANSYS Lumerical software, employing a Gaussian source of 5 μ m diameter to irradiate the PCM with a power of 300 mW and a duration of 300 ns. Initially, a simulation was performed using the FDTD solver to calculate the optical power absorbed per unit volume due to the material absorption, which is then integrated over the entire volume of the GeTe to obtain the total absorption. This parameter is subsequently transferred to the HEAT solver to import the optical absorption data and use it as a heat source in the simulation. The graphs in Figure 4 illustrate the temperature evolution of GeTe as a function of its width and length, recorded on the monitors (represented as a dashed line in Figure 2). Notably, the critical temperature required for GeTe amorphization is exceeded in the

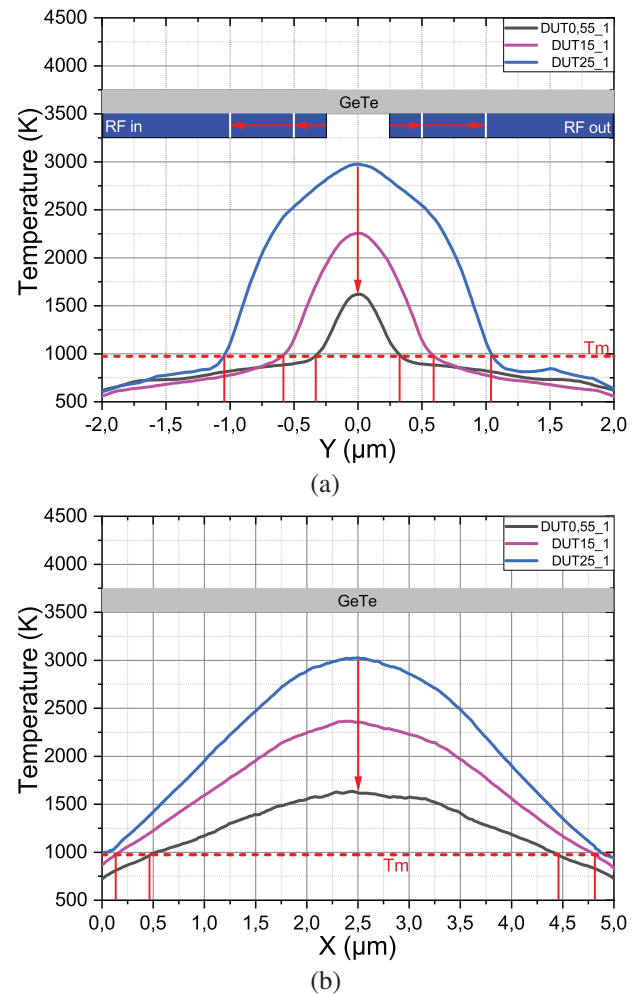


Figure 5. Temperature variation for different RF gaps depending on: (a) PCM length and (b) PCM width.

Table 1. Different devices of GeTe studied in this work

Name	PCM length	PCM width	PCM thick	Optical fiber
DUT0.55_1	0.5 μ m	5 μ m	100 nm	Hi 780 $\varnothing$ 5 μ m
DUT15_1	1 μ m	5 μ m	100 nm	Hi 780 $\varnothing$ 5 μ m
DUT25_1	2 μ m	5 μ m	100 nm	Hi 780 $\varnothing$ 5 μ m
DUT15_2	1 μ m	5 μ m	200 nm	Hi 780 $\varnothing$ 5 μ m
DUT25_2	2 μ m	5 μ m	200 nm	Hi 780 $\varnothing$ 5 μ m

Note: Hi780 : high index fiber with operating λ of 780 nm.

structure with a 5 μ m width and a 2 μ m RF gap, as indicated by the temperature surpassing the melting threshold over the entire width of our GeTe.

Using the same structure and energy input, we varied the RF gap. The results, illustrated in Figure 5(a) and (b), where the nomenclature stands for Device Under Test with an RF gap varying between 0.5 and 2 μ m in X, a width of 5 μ m in Y and a thickness of 100 nm in Z, indicate that decreasing the RF gap significantly enhances heat dissipation towards the metallic contacts (RF pads). This is evident from the 1500 K temperature drop observed at the center of the GeTe between structures DUT25_1 and DUT0.55_1. The increased thermal dissipation complicates

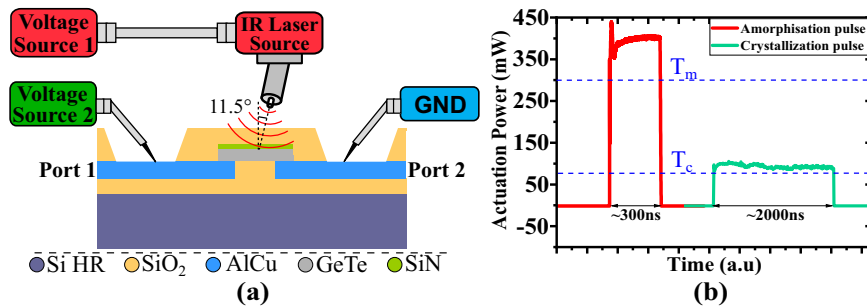


Figure 6. (a) Experimental setup for measurement; (b) Example of laser pulses used in this work.

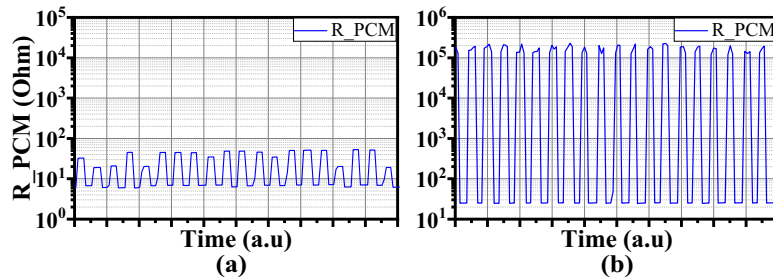


Figure 7. Resistance computation results for different GeTe lengths: (a) For 0.5 μm of length; (b) For 2 μm of length.

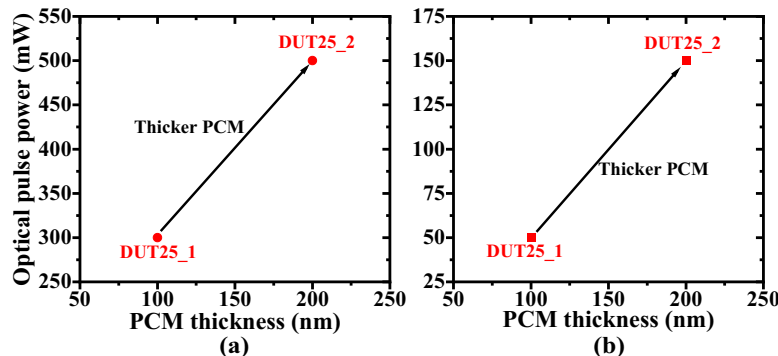


Figure 8. Optical pulse power to amorphize (a) and crystallize (b) PCMs of various thicknesses of 100 and 200 nm, respectively.

the switching of smaller gaps, as parts of the PCM do not reach the melting temperature due to the proximity of the metal, which has a much higher thermal conductivity than GeTe. This highlights the difficulty of achieving a complete phase change in structures with smaller RF gaps, as will be demonstrated in the subsequent measurement section.

Measurement results

Optical command optimization

Upon completion of the optical and thermal simulations, the RF switches were tested using GeTe samples of various lengths and thicknesses, as summarized in Table 1. Optical activation was performed by direct irradiation using an optical fiber connected to a pulsed laser source (915LD-1-3-2) operating at a wavelength of 915 nm. The fiber is positioned 10 μm above the PCM element, with an incidence angle of 11.5° relative to the surface normal, as illustrated in Figure 6(a). A Hi780 optical fiber with a diameter of 5 μm is used to match the width of the 5 μm switches. Micro-positioners from a standard Cascade automatic probing station enable precise alignment of the fiber with the PCM at the wafer level, allowing the delivery of a Gaussian light energy that is tunable and highly localized to the PCM switches. These light pulses are absorbed by the material, causing localized heating within the

PCM. The resulting temperature increase, followed by cooling, induces a phase transition in the PCM, thereby altering its optical and electrical properties, especially its resistance. The devices are actuated with two types of optical pulses, as shown in Figure 6(b). A powerful, short pulse (~ 300 ns) for amorphization and a less powerful, longer pulse (~ 2000 ns) for crystallization. To extract the DC resistance of the PCM, a two-probe measurement technique is used by applying a 10 mV pulse with a Keithley B2902A source, after each optical activation. Figure 7 illustrates the variation in DC resistance between the amorphous and crystalline states after each pulse. It is observed that with a switch having a 0.5 μm RF gap (Figure 7(a)), the switching of GeTe fails, with an $R_{\text{OFF}}/R_{\text{ON}}$ ratio below 20. This failure is attributed to the partial amorphization of GeTe, which does not reach its melting temperature uniformly because of the Gaussian intensity distribution and the high heat dissipation from the metal (RF pads) closer to the center of the GeTe. Conversely, with a switch having a 2 μm RF gap, the GeTe switches correctly, achieving an $R_{\text{OFF}}/R_{\text{ON}}$ ratio exceeding 10^4 , as shown in Figure 7(b). These results confirm the findings from our combined optical and thermal simulations. Figure 8 illustrates the minimum power required to amorphize the DUT25_1 and DUT25_2 devices when optical pulses with increments of 50 mW are applied until switching is achieved. It is noteworthy that the switches are illuminated by a 5 μm fiber with an almost Gaussian profile, indicating that the light is localized at the center.

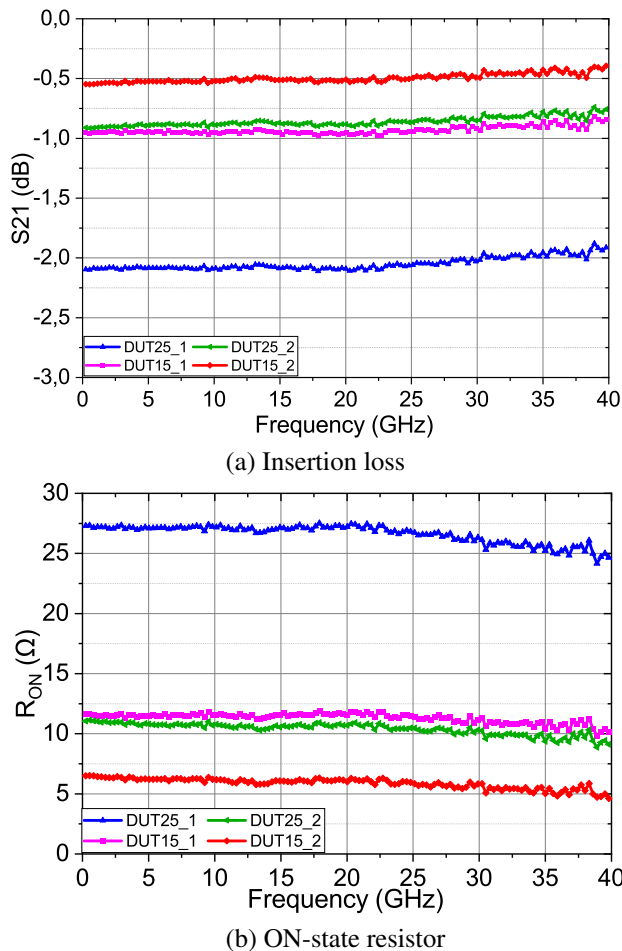


Figure 9. Crystalline GeTe RF performances for different sizes.

To ensure sufficient temperature at the edges of the sample, the comparison between DUT25_1 and DUT25_2 highlights the effect of thickness. It is evident that the required switching powers are not proportional to the GeTe volume. Specifically, DUT25_2, which has twice the volume of GeTe compared to DUT25_1, requires more than twice the power for switching (Figure 8(a)). This can be explained by the fact that optical absorption follows an exponential decay with thickness. Figure 8(b) leads to similar conclusions for the crystallization process.

RF measurement

The RF performance of optically actuated SPST PCM switches was evaluated through S-parameter measurements performed on all devices using a Keysight PNAX vector network analyzer with two ports. Calibration was conducted over the frequency range of 40 MHz to 40 GHz. Initially, the analysis focused on switches in the ON state (optically crystallized PCM). Figure 9(a) illustrates the measured insertion loss after de-embedding in this state for switches of various lengths and thicknesses. The observations indicate that the insertion loss is halved by either doubling the PCM cross-sectional area or halving its length. These results are subsequently used to extract the on-state resistance (R_{ON}) up to 40 GHz, as shown in Figure 9(b). This figure highlights the reduction in resistance from 12 Ω (DUT15_1) to 6 Ω (DUT15_2)

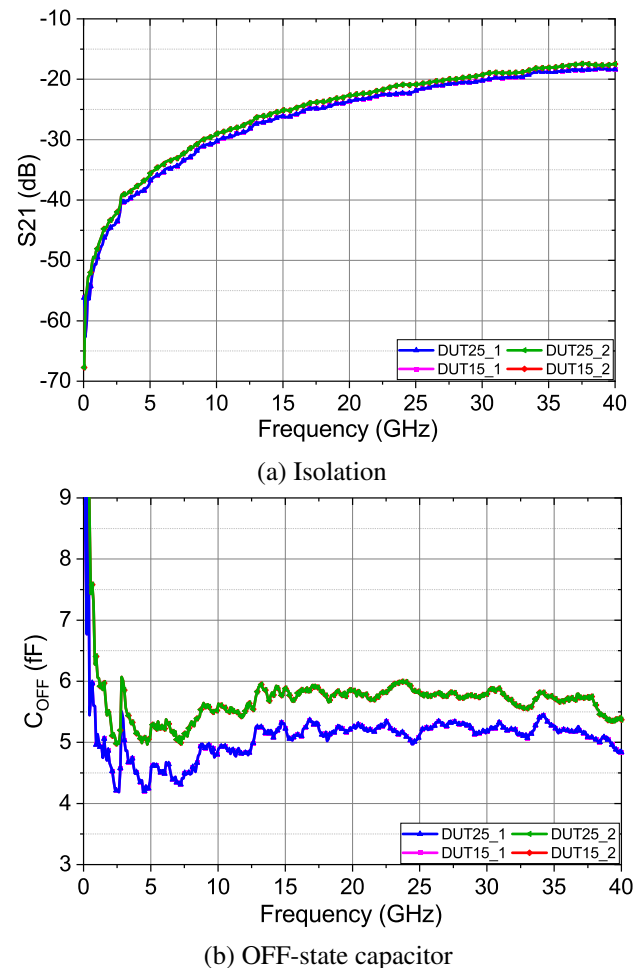


Figure 10. Amorphous GeTe RF performances for different sizes.

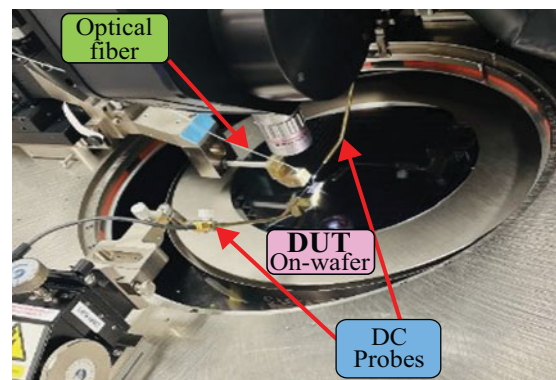


Figure 11. Photograph of the optical measurement test bench on the wafer.

when the cross-sectional area of the PCM switch is doubled while maintaining the same length of 1 μm . This outcome corroborates the theoretical calculations presented in Table 2, emphasizing the advantages of thicker PCM layers. Figure 10(a) displays the S-parameters of the switches in the OFF state. In this scenario, the SPST switches are optically actuated by a short pulse to achieve the amorphous state, characterized by an excellent isolation exceeding 19 dB across the entire frequency range up to 40 GHz. The OFF-state capacitances are extracted and presented in Figure 10(b).

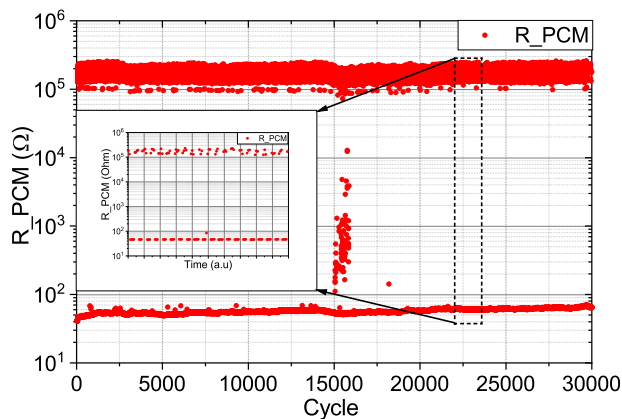


Figure 12. Switching cycles measurement.

As anticipated, the thicker devices exhibit slightly higher capacitances (but no more than 0.8 fF of additional capacitance) due to the doubled coupling area. This results in a figure of merit ($\text{FoM} = R_{\text{ON}} \cdot C_{\text{OFF}}$) of 31.5 fs with a cutoff frequency of 5 THz for device DUT15_2, compared to 57.9 fs and 2.75 THz for device DUT15_1.

Cycling

After determining the conditions for optical actuation, the switching cycle is studied on a switch with a length of 4 μm , a width of 5 μm and a thickness of 100 nm. Using the configuration of Figure 11, applying laser pulses (one for amorphization, then one for crystallization) with a repetition rate of 4 Hz as in Figure 6(b), the DC resistance was measured over several hours for both states after each transition, as shown in Figure 12. DC resistance is plotted as a function of the number of cycles, demonstrating that 30k cycles are achieved with an $R_{\text{OFF}}/R_{\text{ON}}$ ratio greater than 10^4 . Moreover, the crystalline and amorphous DC resistance values remain stable throughout the measurements, indicating that the cycling process is conducted effectively. However, some OFF-state DC resistance

values are observed below $10^5 \Omega$, suggesting that GeTe did not fully amorphize as expected. Specifically, 30 failed amorphization pulses are counted. These PCM amorphization failures are due to micro-vibrations of the positioners caused by external elements that temporally misaligning the fiber, leading to partial amorphization of the PCM. Nevertheless, this did not cause any degradation to our switch. Despite these measurement limitations, this represents the first demonstration of an optically actuated PCM switch with such a high number of cycles, whereas only a dozen of cycles have been achieved so far [10, 12].

Discussion

Table 2 compares the endurance of the switching cycles and the FoM of our work with those published in the literature, considering both indirect electrical and optical excitations at a wavelength of 248 nm, which is incompatible with current integrated photonic technology. Regarding the optical excitation, significant performance can be observed in terms of switching cycles. As previously mentioned, this is the first time such a high number of cycles has been achieved with optical excitation, compared to [10, 12]. A FoM of 31.5 fs is obtained, which is comparable to the state of the art, with only 5.5 fs more than in [11], despite their GeTe being ten times larger (even if 3 times longer, this goes in favor of low R_{ON}) and having twice the conductivity of ours. Furthermore, the optical wavelength of 248 nm used in the literature presents particular challenges for integrated photonic, such as material selection, complexity, and manufacturing costs, thus limiting their development, which is not the case for 915 nm. In terms of comparison with electrical excitation, we achieve fewer cycles than [14, 17, 18]. This is due to external factors such as fiber misalignment after hours of cycling, as we did not observe any degradation of our PCM. Comparing the FoM, we can see that our devices present a FoM comparable to the state of the art and also better than our devices with electrical excitation having the same dimensions and GeTe properties, improving from 45.5 to 36 fs for DUT19_1. This can be explained by the excellent control of the crystalline region during optical excitation, as well as the absence of any heating element that induces parasitic capacitances, making C_{OFF} lower.

Table 2. Performance comparison with the state-of-the-art

Ref	Actuation	PCM size	PCM thick	$R_{\text{ON_meas}} (\Omega)$	$R_{\text{ON_cal}} (\Omega)$	$C_{\text{OFF}} (\text{fF})$	FoM meas (fs)	Cycling
[10]	Optical	$2 \times 5 \mu\text{m}^2$	100 nm	23	–	13.79	–	10
[11]	Optical 248nm	$3 \times 50 \mu\text{m}^2$	200 nm	3	–	8.8	26	–
[12]	Optical 248nm	$1.5 \times 0.5 \text{mm}^2$	500 nm	–	–	–	–	10
[6]	Electrical	$0.9 \times 30 \mu\text{m}^2$	– nm	0.9	–	14.1	12.7	–
[15]	Electrical	$2 \times 10 \mu\text{m}^2$	75 nm	4.5	–	35	149	–
[16]	Electrical	$5 \times 44 \mu\text{m}^2$	150 nm	11.5	–	3.3	37	300k
[17]	Electrical	$0.6 \times 12 \mu\text{m}^2$	250 nm	3.9	–	10.2	39.8	–
[18]	Electrical	$3 \times - \mu\text{m}^2$	– nm	4.5	–	5.6	25.2	10M
[19]	Electrical	$- \times 10 \mu\text{m}^2$	– nm	2.3	–	2.7	6.2	1B
This work	Electrical	$1 \times 9 \mu\text{m}^2$	100 nm	6.5	3.8^a	7	45.5	–
This work	Optical 915nm	$1 \times 9 \mu\text{m}^2$	100 nm	6.11	3.8^a	5.9	36	$\approx 30\text{k}$
This work	Optical 915nm	$1 \times 5 \mu\text{m}^2$	200 nm	6	3.5^a	5.25	31.5	$\approx 30\text{k}$

^aCalculated on the basis of the conductivity of [10], which equals to $2.9 \cdot 10^5 \text{ S/m}$.

Conclusion

The optical actuation of an RF switch based on GeTe with various geometries has been studied at a wavelength of 915 nm. The switching conditions for the ON and OFF states have been determined. A cycling test demonstrated a switching capability of at least 30k cycles, representing the highest optical switching rate of PCM switches reported so far. Small signal measurements show a lower FoM for thicker and shorter switches due to their low on-state resistance, confirming the advantage of thicker PCMs under direct optical irradiation. This work paves the way for the realization of high-performance optically actuated RF switches, fully leveraging CMOS-integrated photonics, which will enhance switch performance by eliminating external factors that may affect their performance. Generally, integrated photonics at 915 nm relies on a standard process using silicon nitride optical waveguides [13, 14, 20]. In this context, the co-integration of RF PCM switches and passive integrated photonics represents a significant advancement in optical actuation [21]. Ultimately, the realization of a fully integrated optical actuation circuit can be envisaged using hybrid integrated lasers [22].

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Competing interests. The author(s) declare none.

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